

# Orthogonal and Symplectic Integrals via modulated $2j - k$ bi-orthogonal polynomial systems

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- ① characteristic polynomials  $\Lambda_{A_N}(z)$  of a random matrix  $A_N$  can model  $L$ -functions
- ② Large  $N$  asymptotics of such characteristic polynomials
- ③ Determinants of moments with a modulated or slanted structure - not Toeplitz
- ④ Integrable nature of these particular determinants
- ⑤ Practical tools for computing and studying coefficients

Katz & Sarnak provide evidence of 4 families of  $L$ -functions with the following symmetry types:

- unitary  $U(N)$ ,
- unitary symplectic  $USp(2N)$ ,
- even orthogonal  $O^+(2N)$  and
- odd orthogonal  $O^-(2N)$

$L$ -functions are modelled by characteristic polynomials  $\Lambda_A(z)$  of a random matrix  $A$ .

[Conrey, Rubinstein & Snaith 2006] developed an efficient method to compute such averages in the unitary case:

$$\int_{U(N)} |\Lambda'_A(1)|^{2m} dA = b_m N^{m^2+2m} + O(N^{m^2+2m-1})$$

Explicitly, their coefficient is as follows:

$$b_m = (-1)^m \sum_{h=0}^m \binom{m}{h} \left( \frac{d}{dt} \right)^{h+m} \left( e^{-t} t^{-m^2/2} \det_{m \times m} (I_{m+j-k}(2\sqrt{t})) \right) \Big|_{t=0}$$

where  $I_\ell(x)$  is the modified Bessel function of the first kind.

[Forrester & W 2006] these determinants of  $I$ -Bessel functions can be given in terms of solutions to Painlevé III' differential equations

One then defines

$$\tau_m(t) := 2^{-m(m-1)} t^{-m^2/2} \det_{m \times m} (I_{m+j-k}(2\sqrt{t}))$$

[Forrester & W 2002] knew that this  $\tau_m(t)$  is in fact the Okamoto  $\tau$ -function associated with the Painlevé III'  $\sigma$ -form:

$$(t\sigma'')^2 + \sigma'(4\sigma' - 1)(\sigma - t\sigma') - \frac{1}{4}m^2 = 0$$

This nonlinear second order differential equation has a solution with certain boundary data given in terms of  $\tau_m(t)$  by the formula

$$\sigma_{III,m}(t) = -t \frac{d}{dt} \log (e^{-t/4} t^{m^2} \tau_m(\frac{1}{4}t))$$

Specifying boundary conditions, one can quickly compute  $\sigma_{III,m}(t)$  from the differential equation and recover  $\tau_m(t)$  via the equation

$$\tau_m(t) = \exp \left( - \int_0^{4t} \frac{ds}{s} (\sigma_{III,m}(s) + m^2 - \frac{1}{4}s) \right).$$

This expression allows a much faster computation of the constants  $b_m$  than earlier methods.

## The Other 3 Cases

[Ali Altuğ, Bettin, Petrow, Rishikesh & Whitehead 2014] compute

$$M_m(G(2N), s) := \int_{G(2N)} \left( \Lambda_A^{(s)}(1) \right)^m dA$$

where  $G$  denotes  $\mathrm{USp}$ ,  $\mathrm{SO}$ , or  $\mathrm{O}^-$ , and  $dA$  is the Haar measure on  $G$ .

As  $N \rightarrow \infty$ , this models the  $m$ th moment of  $L^{(s)}(1/2)$  in a family of symmetry type  $G$

$$\int_{G(2N)} \left( \Lambda_A^{(s)}(1) \right)^m dA = b_m N^{\ell_m} + O(N^{\ell_m - 1})$$

$s$  is the differentiation order

$b_m$  is a geometrical constant

$\ell_m$  is an exponent dependent upon  $G$

## Their Results

For  $u \in \mathbb{C}$  and  $l \in \mathbb{Z}$  let

$$w_l(u) = \oint_{|z|=1} \frac{e^{z+uz^{-2}}}{z^{l+1}} \frac{dz}{2\pi i} = \frac{1}{\Gamma(l+1)} {}_0F_2 \left( ; \frac{1}{2}l + 1, \frac{1}{2}(l+1); \frac{1}{4}u \right)$$

The role of the  $\tau$ -function is now played by

$$D_n^{(r)}(u) := \det \left( w_{2j-k+r}(u) \right)_{0 \leq j, k \leq n-1}, \quad n \geq 0$$

**USp:**  $M_m(\mathrm{USp}(2N), 2) = b_m(\mathrm{USp}(2N), 2) \cdot (2N)^{\frac{m^2+5m}{2}} + O(N^{\frac{m^2+3m}{2}})$  where

$$b_m(\mathrm{USp}(2N), 2) = 2^{-\frac{m^2+5m}{2}} \frac{d^m}{du^m} \left( e^u D_m^{(1)}(2u) \right) \Big|_{u=0}$$

**O<sup>+</sup>:**  $M_m(\mathrm{O}^+(2N), 2) = b_m(\mathrm{O}^+(2N), 2) \cdot (2N)^{\frac{m^2+3m}{2}} + O(N^{\frac{m^2+m}{2}})$  where

$$b_m(\mathrm{O}^+(2N), 2) = 2^{-\frac{m^2+m}{2}} \frac{d^m}{du^m} \left( e^u D_m^{(0)}(2u) \right) \Big|_{u=0}$$

**O<sup>-</sup>:**  $M_m(\mathrm{O}^-(2N), 3) = b_m(\mathrm{O}^-(2N), 3) \cdot (2N)^{\frac{m^2+5m}{2}} + O(N^{\frac{m^2+3m}{2}})$  where

$$b_m(\mathrm{O}^-(2N), 3) = 3 \cdot 2^{-\frac{m^2+3m}{2}} \frac{d^m}{du^m} \left( e^u D_m^{(1)}(2u) \right) \Big|_{u=0}$$

Consider an arbitrary sequence of moments  $\{w_l\}_{l \in \mathbb{Z}}$ .

Let  $D_n^{(r)}$  denote the  $n \times n$  matrices of  $2j - k$  structure and denote their determinants by  $D_n$ :

$$D_n^{(r)} := \det \begin{pmatrix} w_r & w_{r-1} & \cdots & w_{r-n+1} \\ w_{r+2} & w_{r+1} & \cdots & w_{r-n+3} \\ \vdots & \vdots & \vdots & \vdots \\ w_{r+2n-2} & w_{r+2n-3} & \cdots & w_{r+n-1} \end{pmatrix} \equiv \det_{0 \leq j, k \leq n-1} (w_{r+2j-k})$$

Joint density function has the form

$$\prod_{1 \leq j < k \leq n} (\zeta_k - \zeta_j)(\zeta_k^{-2} - \zeta_j^{-2}), \quad \zeta_l := e^{i\theta_l} \in \mathbb{T}, \quad \theta_l \in (-\pi, \pi]$$

which is distinct from modulus thereof, i.e.

$$\prod_{1 \leq j < k \leq n} |\zeta_k - \zeta_j|^2 |\zeta_k + \zeta_j|$$

Recall

$$D_n^{(r)} := \det \begin{pmatrix} w_r & w_{r-1} & \cdots & w_{r-n+1} \\ w_{r+2} & w_{r+1} & \cdots & w_{r-n+3} \\ \vdots & \vdots & \vdots & \vdots \\ w_{r+2n-2} & w_{r+2n-3} & \cdots & w_{r+n-1} \end{pmatrix} \equiv \det_{0 \leq j, k \leq n-1} (w_{r+2j-k})$$

### Definition

For each offset value  $r \in \mathbb{Z}$ , define the  $2j - k$  sequences of monic polynomials  $\{P_n(z; r)\}_{n=0}^{\infty}$  and  $\{Q_n(z; r)\}_{n=0}^{\infty}$ ,  $\deg P_n(z; r) = \deg Q_m(z; r) = n$ , satisfying the *bi-orthogonality* condition:

$$\int_{\mathbb{T}} P_m(\zeta; r) Q_n(\zeta^{-2}; r) \zeta^{-r} \frac{d\mu(\zeta)}{2\pi i \zeta} = h_n^{(r)} \delta_{mn}, \quad m, n \in \mathbb{N} \cup \{0\},$$

where  $h_n^{(r)}$  is the *norms* of the polynomials squared and  $d\mu(\zeta) \equiv w(\zeta)d\zeta$  for some weight function  $w(z)$ .

If  $D_n^{(r)} \neq 0$ , the polynomials  $P_n(z; r)$  and  $Q_m(z; r)$  exist and are uniquely given by

$$P_n(z; r) = \frac{1}{D_n^{(r)}} \det \begin{pmatrix} w_r & w_{r-1} & \cdots & w_{r-n} \\ w_{r+2} & w_{r+1} & \cdots & w_{r-n+2} \\ \vdots & \vdots & \vdots & \vdots \\ w_{r+2n-2} & w_{r+2n-3} & \cdots & w_{r+n-2} \\ 1 & z & \cdots & z^n \end{pmatrix}$$

Observe  $h_n^{(r)} = D_{n+1}^{(r)}/D_n^{(r)}$

Analogues of the Szegö recurrences are the third order  $n$ -recurrences

$$P_{n+3}(z; r) + \left( \frac{h_{n+2}^{(r-1)}}{h_{n+2}^{(r)}} + \frac{h_{n+1}^{(r-2)}}{h_{n+1}^{(r-1)}} \right) P_{n+2}(z; r) + \left( \frac{h_{n+1}^{(r-2)}}{h_{n+1}^{(r)}} - z^2 \right) P_{n+1}(z; r) - \frac{h_{n+1}^{(r-2)}}{h_n^{(r)}} z^2 P_n(z; r) = 0$$

and for  $Q_n^*(z; r) := z^n Q_n(z^{-1}; r)$

$$Q_{n+3}^*(z; r) - \left( 1 - \frac{h_{n+2}^{(r+2)}}{h_{n+2}^{(r)}} z \right) Q_{n+2}^*(z; r) - \left( \frac{h_{n+2}^{(r+1)}}{h_{n+1}^{(r)}} + \frac{h_{n+2}^{(r+2)}}{h_{n+1}^{(r+1)}} \right) z Q_{n+1}^*(z; r) - \frac{h_{n+2}^{(r+2)}}{h_n^{(r)}} z^2 Q_n^*(z; r) = 0$$

Third order pure-offset  $r$ -recurrence relations for the  $2j - k$  polynomials are given by

$$\frac{D_n^{(r+3)} D_{n+1}^{(r)}}{D_{n+1}^{(r+1)} D_n^{(r+2)}} P_n(z; r+3) - z P_n(z; r+2) - \frac{h_n^{(r)}}{h_n^{(r+1)}} P_n(z; r+1) + z P_n(z; r) = 0$$

and

$$Q_n^*(z; r+3) - Q_n^*(z; r+2) - \frac{h_n^{(r+3)}}{h_n^{(r+1)}} z Q_n^*(z; r+1) + \frac{D_n^{(r)} D_{n+1}^{(r+3)}}{D_{n+1}^{(r+1)} D_n^{(r+2)}} z Q_n^*(z; r) = 0$$

## Pure & Mixed Recurrence Relations

Let

$$P_n(z; r) = z^n + \sum_{j=0}^{n-1} p_{n,j}^{(r)} z^{n-j}, \quad Q_n(z; r) = z^n + \sum_{j=0}^{n-1} q_{n,j}^{(r)} z^{n-j}$$

Two pure  $n$ -recurrence relations for the sub-leading coefficients

$$p_{n+2,1}^{(r)} - p_{n,1}^{(r)} = -\frac{h_{n+1}^{(r-1)}}{h_{n+1}^{(r)}} - \frac{h_n^{(r-2)}}{h_n^{(r-1)}} + \frac{h_n^{(r-2)}}{h_{n-1}^{(r)}}, \quad q_{n+1,1}^{(r)} - q_{n,1}^{(r)} = -\frac{h_n^{(r+2)}}{h_n^{(r)}} + \frac{h_n^{(r+1)}}{h_{n-1}^{(r)}} + \frac{h_n^{(r+2)}}{h_{n-1}^{(r+1)}}$$

Two pure  $r$ -recurrences

$$p_{n,1}^{(r+2)} - p_{n,1}^{(r)} = -\frac{h_n^{(r)}}{h_{n-1}^{(r+2)}}, \quad q_{n,1}^{(r+1)} - q_{n,1}^{(r)} = \frac{h_n^{(r+1)}}{h_{n-1}^{(r)}}$$

and in total there are 4 mixed versions such as

$$q_{n+1,1}^{(r)} - q_{n,1}^{(r+2)} = -\frac{h_n^{(r+2)}}{h_n^{(r)}}$$

Reproducing kernel is defined by

$$K_n(z, \zeta; r) := \sum_{j=0}^n \frac{1}{h_j^{(r)}} Q_j(z; r) P_j(\zeta; r)$$

The Christoffel-Darboux identity for the "2j - k" system can be written as

$$\begin{aligned} K_n(z_2, z_1; r) &= \frac{1}{z_2 z_1^2 - 1} \left\{ \frac{1}{h_n^{(r+2)}} P_n(z_1; r+2) Q_{n+1}(z_2; r) + \frac{1}{h_n^{(r)}} z_2 P_{n+2}(z_1; r+2) Q_n(z_2; r) \right. \\ &\quad \left. + \frac{1}{h_{n+1}^{(r+2)}} P_{n+1}(z_1; r+2) \left[ Q_{n+2}(z_2; r) - (z_2 + q_{n+2,1}^{(r)} - q_{n+1,1}^{(r+2)}) Q_{n+1}(z_2; r) \right] \right\} \end{aligned}$$

Recall

$$w(z; u) = e^{z + uz^{-2}}$$

Pearson-type relation  $w' := \partial_z w$

$$\frac{w'}{w} = (1 - 2uz^{-3})$$

Spectral derivatives

$$zP'_n(z; r) = P_{n+1}(z; r) + \left( n + p_{n,1}^{(r)} - p_{n+1,1}^{(r)} - z \right) P_n(z; r) + 2u \frac{h_n^{(r)}}{h_{n-1}^{(r)}} P_{n-1}(z; r)$$

and

$$zQ'_n(z; r) = uQ_{n+1}(z; r) + \left[ n + u \left( q_{n,1}^{(r)} - q_{n+1,1}^{(r)} - z \right) \right] Q_n(z; r) + \frac{1}{2} \frac{h_n^{(r)}}{h_{n-1}^{(r)}} Q_{n-1}(z; r)$$

Considering

$$\int_{\mathbb{T}} P_n(\zeta; r) Q_m(\zeta^{-2}; r) \zeta^k (1 - 2u\zeta^{-3}) w(\zeta; u) \zeta^{-r} \frac{d\zeta}{2\pi i \zeta}$$

it is possible to deduce a sequence of identities  $k = \pm 1, 3, \dots$  for  $m - n = 0, \pm 1, \dots$

E.g. for  $k = 1, m = n$ :

$$p_{n,1}^{(r)} - 2uq_{n,1}^{(r)} = -nr - \frac{1}{2}n(n-1)$$

Pearson-type relation  $\dot{w} := \partial_u w$

$$\frac{\dot{w}}{w} = z^{-2}$$

Linear differential-difference equations for the weight

$$\partial_u w_r = w_{r+2}$$

Third order linear differential equation

$$4u^2 \partial_u^3 w_r + (4r + 10)u \partial_u^2 w_r + (r + 1)(r + 2)\partial_u w_r - w_r = 0$$

Third order linear difference equation

$$2uw_{r+3} + (r + 1)w_{r+1} - w_r = 0$$

Polynomial deformation derivatives are:

$$\dot{P}_n(z; r) = -\frac{h_n^{(r)}}{h_{n-1}^{(r)}} P_{n-1}(z; r)$$

$$\dot{Q}_n(z; r) = Q_{n+1}(z; r) + Q_n(z; r) \left( q_{n,1}^{(r)} - q_{n+1,1}^{(r)} - z \right)$$

Define the rank 3 system polynomial variables, all with assumed offset  $r$

$$\mathcal{P}_n(z; u) := \begin{pmatrix} P_{n+1}(z; u) \\ P_n(z; u) \\ P_{n-1}(z; u) \end{pmatrix}, \quad Q_n(z; u) := \begin{pmatrix} Q_{n+1}(z; u) \\ Q_n(z; u) \\ Q_{n-1}(z; u) \end{pmatrix}$$

N.B. Other choices of linearly-independent triples are possible.

Transfer operators

$$\mathcal{P}_{n+1} = \mathcal{M}_n \mathcal{P}_n, \quad Q_{n+1} = \mathcal{N}_n Q_n$$

with the transfer matrices given by

$$\mathcal{M}_n(z; r) = z^2 \begin{pmatrix} 0 & 1 & \frac{h_n^{(r-2)}}{h_{n-1}^{(r)}} \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} + \begin{pmatrix} -\frac{h_{n+1}^{(r-1)}}{h_{n+1}^{(r)}} - \frac{h_n^{(r-2)}}{h_n^{(r-1)}} & -\frac{h_n^{(r-2)}}{h_n^{(r)}} & 0 \\ 1 & 0 & 0 \\ 0 & 1 & 0 \end{pmatrix}$$

N.B. Another discrete structure is the offset shift  $r \mapsto r + 1$ .

The spectral derivatives are denoted

$$\partial_z \mathcal{P}_n = \mathcal{A}_n \mathcal{P}_n, \quad \partial_z \mathcal{Q}_n = \mathcal{U}_n \mathcal{Q}_n$$

The spectral matrices are

$$\begin{aligned} \mathcal{A}_n(z; r) &= z \begin{pmatrix} 0 & 1 & \frac{h_n^{(r-2)}}{h_{n-1}^{(r)}} \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} + \begin{pmatrix} -1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & -1 \end{pmatrix} \\ &+ \frac{1}{z} \begin{pmatrix} n+1 + p_{n+1,1}^{(r)} - p_{n+2,1}^{(r)} - \frac{h_{n+1}^{(r-1)}}{h_{n+1}^{(r)}} - \frac{h_n^{(r-2)}}{h_n^{(r-1)}} & 2u \frac{h_{n+1}^{(r)}}{h_n^{(r)}} - \frac{h_n^{(r-2)}}{h_n^{(r)}} & 0 \\ 1 & n + p_{n,1}^{(r)} - p_{n+1,1}^{(r)} & 2u \frac{h_n^{(r)}}{h_{n-1}^{(r)}} \\ 0 & 1 & n-1 + p_{n-1,1}^{(r)} - p_{n,1}^{(r)} - 2u \frac{h_{n-1}^{(r)}}{h_{n-1}^{(r-2)}} \end{pmatrix} \\ &+ \frac{1}{z^3} \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 2u \frac{h_{n-1}^{(r)}}{h_{n-1}^{(r-2)}} & 2u \left( \frac{h_{n-1}^{(r)}}{h_{n-1}^{(r-1)}} + \frac{h_{n-1}^{(r)}}{h_{n-1}^{(r-2)}} \frac{h_n^{(r-1)}}{h_n^{(r)}} \right) & 2u \end{pmatrix} \end{aligned}$$

N.B. Two irregular singularities at  $z = 0, \infty$  with Poincaré ranks 2

The deformation derivatives are denoted by

$$\partial_u \mathcal{P}_n = \mathcal{B}_n \mathcal{P}_n, \quad \partial_u \mathcal{Q}_n = \mathcal{V}_n \mathcal{Q}_n$$

The deformation matrices are

$$\mathcal{B}_n(z; r) = \frac{1}{z^2} \begin{pmatrix} 0 & 0 \\ 0 & 0 \\ -\frac{h_n^{(r)}}{h_{n-1}^{(r-2)}} & -\frac{h_{n-1}^{(r)}}{h_{n-1}^{(r-2)}} \frac{h_n^{(r-1)}}{h_n^{(r)}} - \frac{h_n^{(r)}}{h_{n-1}^{(r-1)}} \end{pmatrix} + \begin{pmatrix} 0 & -\frac{h_{n+1}^{(r)}}{h_n^{(r)}} & 0 \\ 0 & 0 & -\frac{h_n^{(r)}}{h_{n-1}^{(r-1)}} \\ 0 & 0 & \frac{h_{n-1}^{(r)}}{h_{n-1}^{(r-2)}} \end{pmatrix}$$

## Compatibility & Integrability relations of the dynamical system

$$\mathcal{A}_{n+1} \cdot \mathcal{M}_n - \mathcal{M}_n \cdot \mathcal{A}_n = \partial_z \mathcal{M}_n$$

$$\mathcal{B}_{n+1} \cdot \mathcal{M}_n - \mathcal{M}_n \cdot \mathcal{B}_n = \partial_u \mathcal{M}_n$$

$$\mathcal{A}_n \cdot \mathcal{B}_n - \mathcal{B}_n \cdot \mathcal{A}_n = \partial_z \mathcal{B}_n - \partial_u \mathcal{A}_n$$

The  $\sigma$ -function

$$\partial_u \log D_n^{(r)} = -q_{n,1}$$

$P - Q$  Linkage

$$p_{n,1}^{(r)} - 2uq_{n,1}^{(r)} = -nr - \frac{1}{2}n(n-1)$$

Closure identities

$$2u \left[ \frac{h_{n+1}^{(r)}}{h_n^{(r-1)}} + \frac{h_{n+1}^{(r-1)}}{h_n^{(r-2)}} \right] + \frac{h_{n+1}^{(r-2)}}{h_n^{(r)}} = 2n + 2$$

First derivatives

$$\partial_u \log h_n^{(r)} = \frac{h_n^{(r)}}{h_n^{(r-2)}} - \frac{h_{n+1}^{(r-1)}}{h_n^{(r-2)}} - \frac{h_{n+1}^{(r)}}{h_n^{(r-1)}}$$

$$u\partial_u q_{n,1}^{(r)} + q_{n,1}^{(r)} + \frac{1}{2} \frac{h_n^{(r)}}{h_{n-1}^{(r)}} = 0$$

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